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Lithium niobate photonic crystal waveguides: Far field and near field characterisation

M.-P. Bernal *, N. Courjal, J. Amet, M. Roussey, C.H. Hou

Institut FEMTO-ST (UMR 6174), Département d'Optique P.M. Duffieux, Université de Franche-Comté, 25030 Besançon Cedex, France

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8 Abstract

9 In this paper, we experimentally investigate photonic crystal waveguides in a X-cut lithium niobate substrate. The transmission
10 response is measured through the ΓM direction of a triangular lattice structure and the results coincide with the theoretical predictions.
11 In addition, a scanning near-field microscope is used in collection mode to map the optical intensity distribution inside the structure putt-
12 ing in evidence the guiding of the light through lines of defects. This study offers perspectives towards lithium niobate tunable photonic
13 crystal devices.

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16 Keywords: Lithium niobate photonic crystal; SNOM; Waveguides

18 Photonic crystals (PCs), also known as photonic band-
19 gap materials, are attractive optical materials for control-
20 ling and manipulating the flow of light. Their structure
21 consists basically on periodic changes of the dielectric con-
22 stant on a length scale comparable to optical wavelengths.
23 Multiple interference between scattered light waves can
24 eventually lead to some frequencies that are not allowed
25 to propagate, giving rise to forbidden and permitted bands,
26 similar to the electronic bandgap in a semiconductor. The
27 band structure depends on the geometry and the material
28 refractive index. Hence, an attractive feature of photonic
29 crystals consists in tuning the substrate refractive index con-
30 trolling therefore the transmission response. With tunable
31 photonic crystals, the path is open towards high density
32 ultra-compact photonic circuits. This perspective has moti-
33 vated various studies on tunable photonic devices [1–5].

34 Among optical tunable materials, the combination of
35 excellent electro-optical, acousto-optical, non-linear optical
36 properties, electro-mechanical (piezoelectric) properties,

chemical and mechanical stability makes lithium niobate
(LN) an attractive host material for application in photonic
crystal devices. Moreover, LN high electro-optical coeffi-
cient and its low optical losses make it very adequate for
optical communication systems.

In our two previous works [6,7], we have shown the fab-
rication by focused ion beam (FIB) milling of a triangular
lattice of nanometric-sized holes with etching depths of
2 μm on an annealed proton exchanged (APE) lithium ni-
obate waveguide [8]. We have also shown both theoretically
as well as experimentally the presence of a photonic band-
gap (PBG) with an extinction ratio lower than -12 dB.

Recently, an alternative fabrication technique that con-
sists of electric poling and subsequent etching has con-
firmed the interest in LN-based nanodevices [8].

In this work, the possibility of guiding the light is exper-
imentally evaluated for photonic crystal waveguides. We
present a far field as well as a scanning near field experi-
mental characterization of a photonic crystal waveguide
fabricated on a lithium niobate waveguide.

As mentioned above, our final objective is the fabrication
of photonic bandgap structures in which transmission can

* Corresponding author. Tel.: +33 3 81 66 64 10; fax: +33 3 81 66 64 23.
E-mail address: maria-pilar.bernal@univ-fcomte.fr (M.-P. Bernal).

59 be tuned by changing the refractive index. We have already
 60 theoretically shown [7] that in the case of a triangular array
 61 of holes, the optimal sensitivity to the refractive index is
 62 obtained when the direction of propagation is ΓM , and
 63 the polarization of the electric field is TE (parallel to the sub-
 64 strate plane and perpendicular to the direction of the holes).
 65 The ΓM propagation direction exhibits the additional
 66 advantage of requiring a lower number of rows to obtain
 67 a photonic gap. Indeed, we have shown in Ref. [7] that the
 68 ΓM direction requires only 15 rows to get a -12 dB extinc-
 69 tion ratio as opposed to the 30 rows that would be necessary
 70 to achieve the same gap in the propagation direction ΓK .
 71 Due to the well-known difficulty to etch lithium niobate, this
 72 property has strongly motivated our choice. In addition,
 73 with such a configuration we have experimentally demon-
 74 strated the existence of a photonic gap.

75 To complete the analysis of this configuration, we have
 76 fabricated two alternative structures, based on the same
 77 array as in our previous work, but with one (PCW1) or
 78 three lines (PCW3) of defects. The aim is to investigate
 79 the possibilities of a tunable guiding of the light through
 80 the crystal. The geometrical parameters are chosen to get
 81 a transmission zone around 1550 nm within the bandgap.

82 The photonic crystals are fabricated on a 0.3 mm thick X-
 83 cut LiNbO_3 wafer. In a first step, an optical gradient index
 84 waveguide is fabricated by annealed proton exchange. This
 85 step is realized through a SiO_2 mask in benzoic acid at
 86 180 °C during 1.5 h. The process is followed by an annealing

87 of the optical waveguide at 333 °C for 9 h. These parameters
 88 are chosen to position the core of the optical mode as close as
 89 possible to the surface (approximately 1.4 μm) while keeping
 90 single mode propagation at 1.55 μm .

91 The photonic crystal structure was fabricated in the cen-
 92 tral region on the optical channel waveguide as shown in
 93 Fig. 1(a). It consists of a triangular lattice of 48×26 circular
 94 holes. The lithium niobate substrate (300 microns thickness)
 95 is metalised with a thin Cr layer (100 nm) to avoid charging
 96 effects. This Cr layer is deposited by electron gun evapora-
 97 tion (Balzer, B510). The sample is g rounded with a conduc-
 98 tive paste before introduction in the FIB vacuum chamber
 99 (10^{-6} Torr). The FIB used is a FEI Dual Beam Strata 235.
 100 Ga^+ ions are emitted with an accelerated voltage of 30 keV
 101 and focused down with electrostatic lenses on the sample
 102 with a probe current of 120 pA. The Gaussian beam shape
 103 spot size is about 20 nm at the sample surface. The etching
 104 time of the structures PCW1 and PCW3 (48×26 triangular
 105 hole lattice, hole diameter = 255 nm, periodicity = 510 nm,
 106 etching depth = 1500 nm) was 20 min each. We would like
 107 to point out that the removal of material by FIB milling is
 108 achieved without the use of a patterned resist mask and
 109 therefore, high-precision complex structures can be directly
 110 fabricated. A FIB image cross-section of the holes is shown
 111 in Fig. 1(b). The angle between the FIB beam and the holes
 112 axis is 52° .

113 In order to couple the conventional TE APE mode
 114 (4 μm size) to the photonic crystal waveguides (approxi-

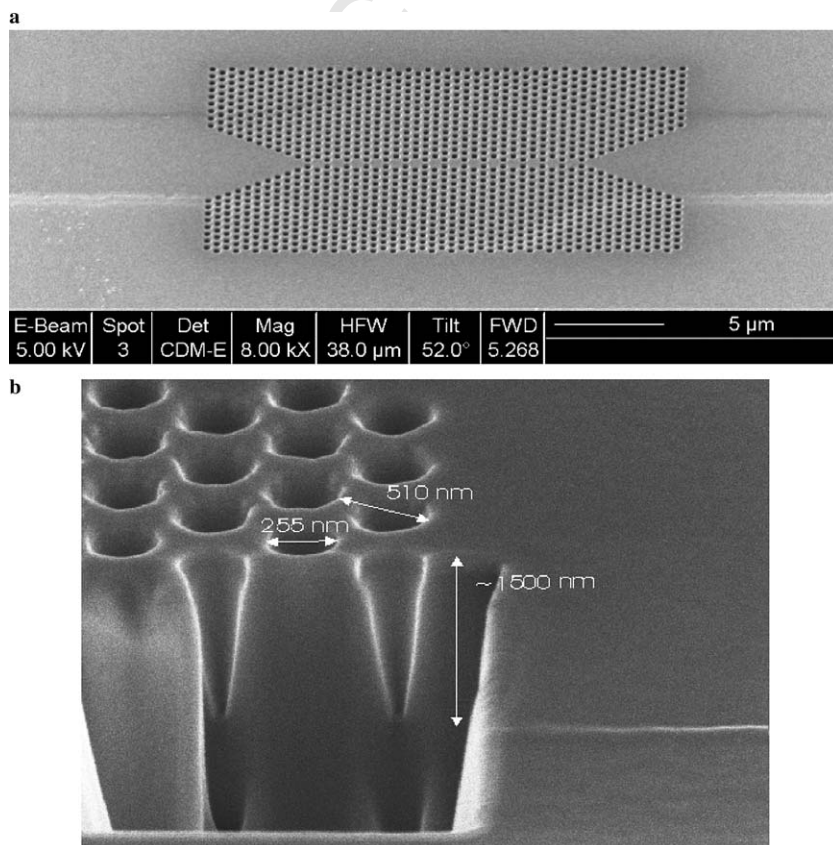


Fig. 1. SEM view of the: (a) triangular lattice PCW1 and (b) its cross-section.

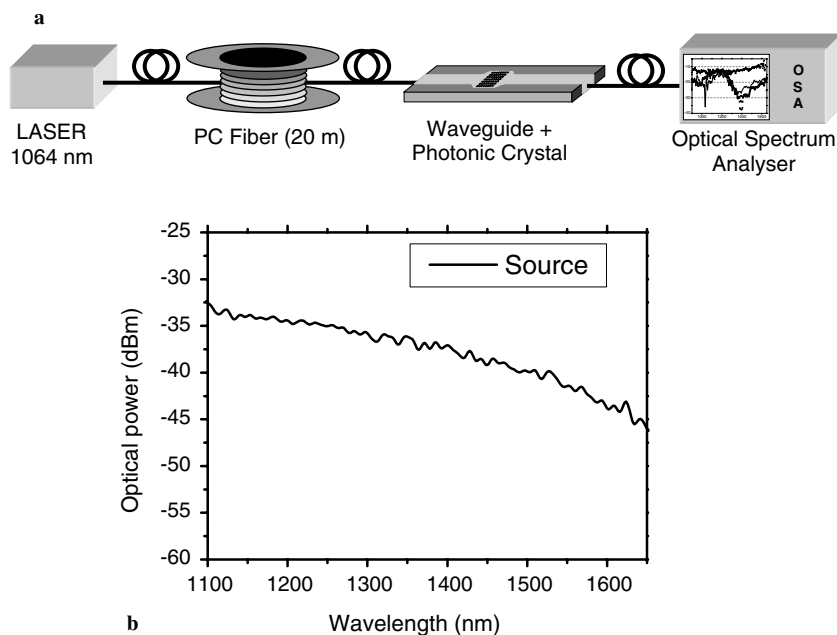


Fig. 2. (a) Experimental setup and (b) super-continuum at the output of the PC fibre.

115 mately 500 nm and 1500 nm wide), the optical mode is
 116 smoothly guided through a photonic tapered structure
 117 (see Fig. 1(a)).

118 The novel structures were first characterized by measur-
 119 ing their far field transmission. The experimental setup is
 120 shown in Fig. 2(a). In order to get a spectrum as flat as pos-
 121 sible on a large range of wavelengths (1000–1700 nm), we

122 use a super-continuum light source. The white light is gen-
 123 erated by a sub-nanosecond microchip laser emitting at
 124 1064 nm with 8 μ J energy per pulse [9]. The laser light is
 125 coupled into a photonic crystal (PC) fibre, which enhances
 126 the nonlinear effects required for the generation of a large
 127 super-continuum. The resulting output spectrum for a
 128 20 m long PC fibre is shown in Fig. 2(b).

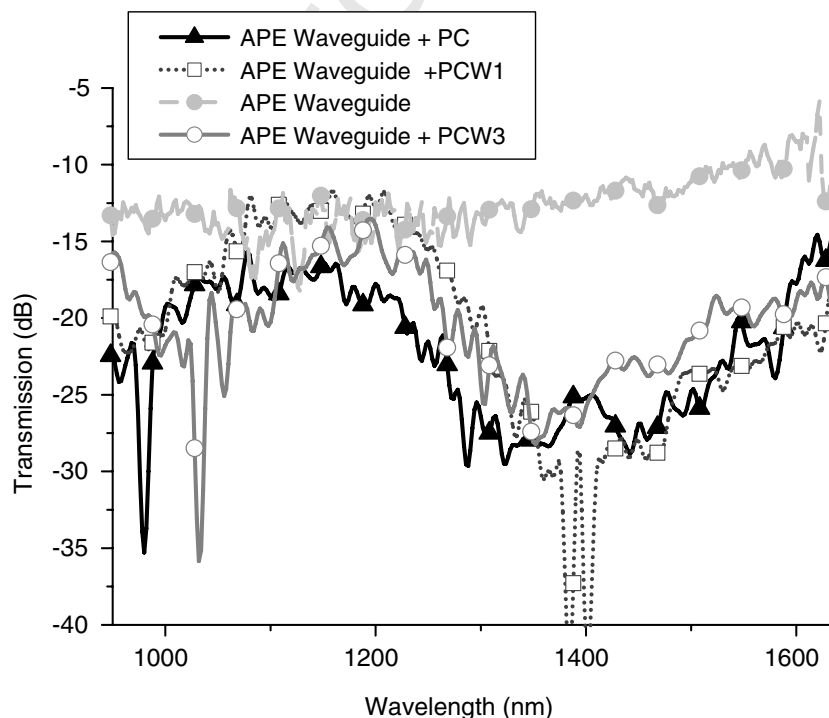


Fig. 3. Experimental transmission through three optical waveguides integrated on the same wafer: Light grey line with solid circle: transmission through an annealed proton exchanged (APE) waveguide. Grey line with empty circles: transmission through a PCW3 integrated on an APE waveguide. Dark grey with empty squares: transmission through a PCW1 integrated on an APE waveguide. Black line with solid triangles: transmission through a PC integrated on an APE waveguide.

129 The optical transmission was measured through the two
 130 photonic crystal waveguides, a photonic crystal without
 131 defect lines, and through a standard optical waveguide,
 132 fabricated on the same wafer and in the same conditions

as described above. The experimental results are shown 133
 in Fig. 3. As it can be seen in the graph, the transmissions 134
 through the photonic structures (filled triangle, empty 135
 square, empty circle) exhibit a gap, which does not appear 136

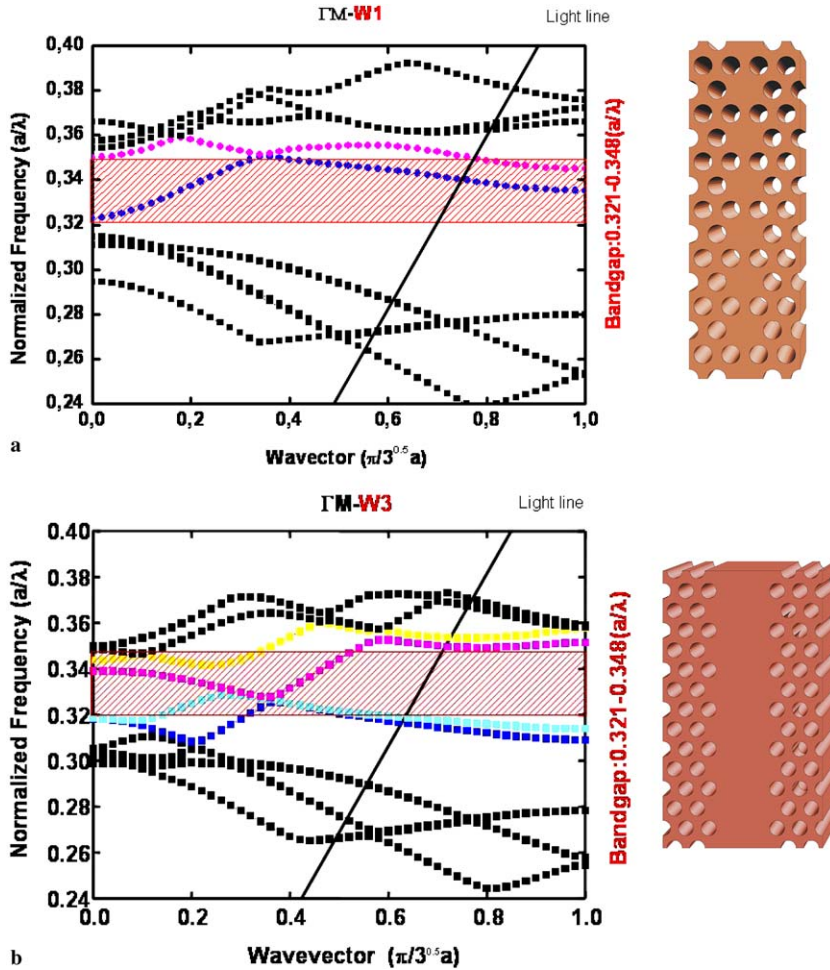


Fig. 4. Projected band diagrams and light line along the ΓM direction for the two photonic crystal waveguides: (a) PCW1 and (b) PCW3.

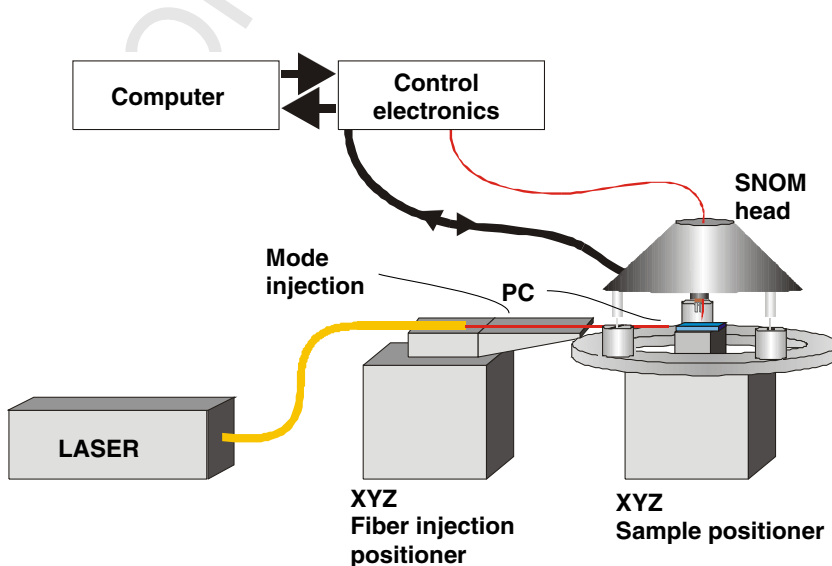


Fig. 5. Experimental set-up of the scanning near-field optical microscope (working in collection mode).

137 in the transmission through the single APE waveguide
 138 (filled circle). In parallel, 2D-numerical simulations per-
 139 formed with a commercial software (BandSOLVE) of the
 140 device without defect lines predict a band gap between
 141 1465 nm and 1589 nm. The experimental gap starts in a
 142 shorter wavelength (approximately around 1300 nm) which
 143 we believe is a consequence of fabrication imperfections
 144 and to the fact that 2D calculated bandgaps are usually
 145 shifted towards shorter wavelengths compared to full 3D
 146 simulations. Fig. 4(a) and (b) show the projected band dia-
 147 grams and the light line for the two photonic crystal wave-
 148 guides. For the PCW1 case, the diagram shows two guided
 149 modes. In the PCW3 case all the modes are radiative
 150 (Fig. 4(b)). Experimentally (Fig. 3), light propagation in
 151 the PC waveguides is observed by an increase in the trans-
 152 mission inside the gap. This increase is twice more impor-
 153 tant for the PCW3 case which may be due to a better
 154 coupling efficiency of the input taper. We have repeated
 155 the measurements five times, and changed the injection
 156 conditions in order to verify the position of the gap and
 157 to optimize the propagation through the photonic crystal
 158 waveguides. The position of the gap did not change in all
 159 the five measurements.

160 For a deeper interpretation of the propagation of the
 161 light through the structures, we have also investigated the
 162 near field behaviour of the light inside the PC waveguides.
 163 The wave fronts of light in the photonic crystal waveguide

164 undergo substantial modulations on length scales that are
 165 much shorter than one wavelength. Therefore, it is impos-
 166 sible to resolve the spatial details of light propagation by
 167 the far field transmission measurement described above.
 168 Although still not systematically used in the photonic cry-
 169 stal community, several groups have shown already very
 170 interesting results in near field characterisation of photonic
 171 crystals [10–17].

172 In the work presented here, the instrument used is a
 173 commercial scanning near-field optical microscope
 174 (SNOM) (NT-MDT SMENA) in collection mode [18].
 175 The near-field optical fibre probe is fabricated by heating
 176 a single mode optical fibre with a CO₂ laser and then pull-
 177 ing it apart with a micropipette puller (Sutter Instrument
 178 Co.) to obtain a sharp taper region with a small end face
 179 (~100 nm). To obtain the SNOM images, one needs to
 180 scatter the evanescent fields on the sample by raster scan-
 181 ning the sub-wavelength probe at a few nanometers from
 182 the surface. A non-optical shear force feedback [19] is used
 183 to keep the probe at a constant distance from the sample
 184 surface. Both signals, the feedback and the optical one,
 185 are simultaneously acquired to construct topographic and
 186 SNOM collection images.

187 Fig. 5 shows the experimental set-up. To image the
 188 transmitted mode through the LN photonic structure,
 189 two different laser sources and two different optical de-
 190 tectors have been utilized. The first acquisition has been per-

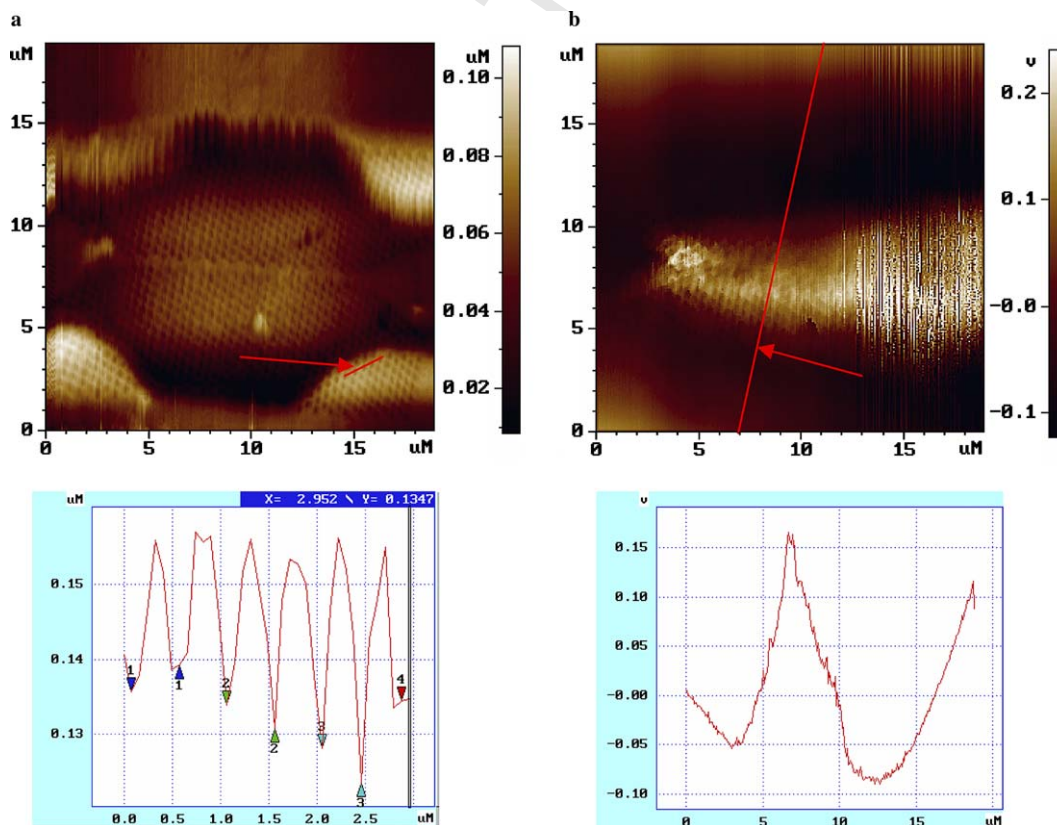


Fig. 6. (a) Topographical image of $20 \times 20 \mu\text{m}$ size of the PCW1 structure. The inset shows a cross-section through 7 holes inside the structure. (b) The simultaneously recorded optical near field of the structure when the coupled wavelength is 810 nm. The inset shows the width of the optical field inside the structure.

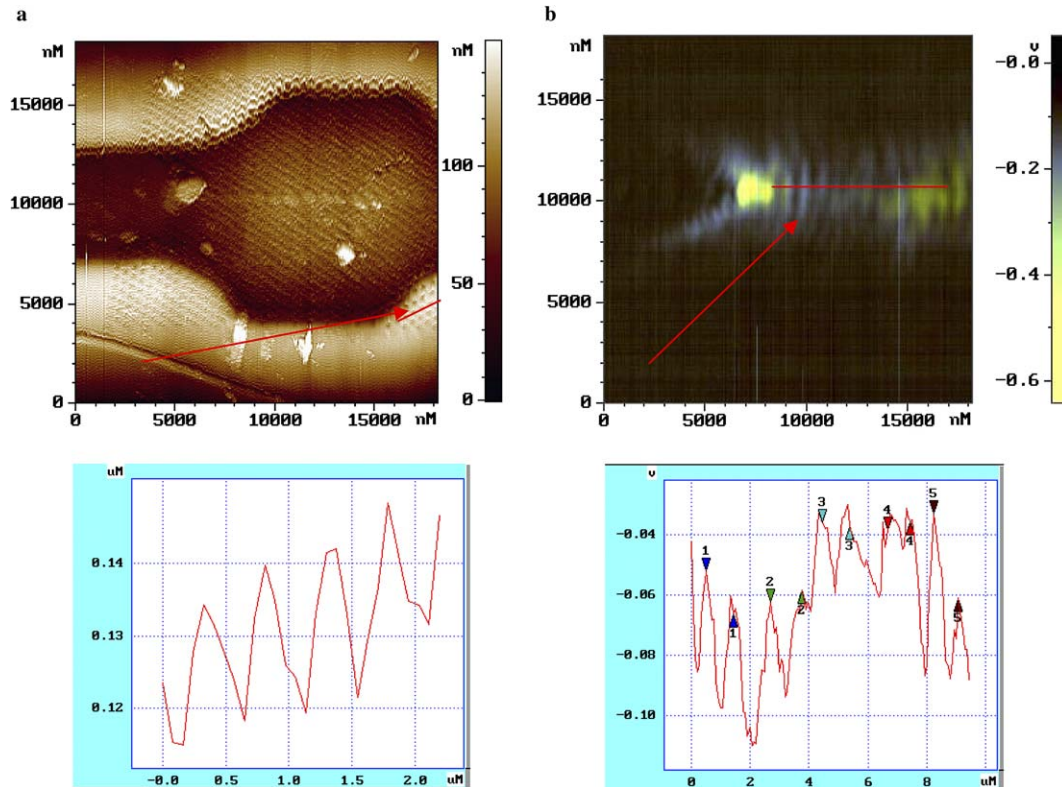


Fig. 7. (a) Topographical image of $20 \times 20 \mu\text{m}$ size of the PCW1 structure. The inset shows a cross-section through 5 holes inside the structure. (b) The simultaneously recorded optical near field of the structure when the coupled wavelength is 1550 nm . The inset shows the cross-section of the optical field along the propagation axes.

191 formed at 810 nm (outside the gap), with a Thorlabs
 192 STFC780 laser and detected with an Oriel 70680 photo-
 193 multiplier. The second acquisition has been realized at
 194 $1.55 \mu\text{m}$ with a distributed feedback laser OKI OL502OON
 195 and an InGaAs detector (Thorlabs D400FC), to character-
 196 ize the propagating region inside the gap.

197 The optical image and topography of the PCW1 at
 198 810 nm is shown in Fig. 6. Fig. 6(a) shows the topography
 199 of the PCW1 structure and a zoom of a small region that
 200 consists of 7 holes. The hole depth measured by the SNOM
 201 tip is of the order of 30 nm which is far from the $1.5 \mu\text{m}$
 202 measured in the SEM image. This is basically due to the
 203 fact that the hole diameter is comparable in size to the
 204 tapered fibre, being difficult for the tip to penetrate inside
 205 the holes. The signal to noise ratio (SNR) is however high
 206 (~ 10). Fig. 6(b) shows the optical image of the light going
 207 through the PCW1 at 810 nm .

208 We have also performed the near field measurements in
 209 a region inside the gap in which an optical mode propa-
 210 gates ($\lambda \sim 1.55 \mu\text{m}$, transmission $\sim -25 \text{ dB}$). The results
 211 are shown in Fig. 7. Fig. 7(a) shows the topography and
 212 Fig. 7(b) the near field image respectively. It is worth men-
 213 tioning that this set of measurements have been performed
 214 with a different SNOM probe than for the case of 810 nm
 215 due to the tip destruction. Again, the topography shows
 216 clearly the photonic structure with a SNR of 10. We can
 217 very well appreciate the tapered beginning and end of the

218 structure and the line of defects. The simultaneously
 219 recorded optical signal is shown in Fig. 7(b). This image
 220 corresponds to an input wavelength of $1.55 \mu\text{m}$. The
 221 recorded signal shows a periodicity of about 800 nm (see
 222 inset in Fig. 7(b)) that corresponds to a Bloch periodicity
 223 of $\sqrt{3} a$ (ΓM direction, triangular lattice). With these
 224 results, we can infer that the step seen in the transmission
 225 response of the PCW1 (Fig. 3) around 1500 nm is due to
 226 the existence of a guiding region.

227 In conclusion, we have shown the first photonic crystal
 228 waveguides fabricated by FIB etching in lithium niobate.
 229 In spite of the optical losses that are mainly due to multi-
 230 mode guiding of the standard APE waveguide and the fab-
 231 rication imperfections caused by the well-known difficulty
 232 of etching lithium niobate, a guiding mode is successfully
 233 observed in the PCW1 structure by near field
 234 characterization.

235 Work is in progress to show experimentally a tunable
 236 lithium photonic device with an optimised photonic crystal
 237 configuration.

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